

# Improved parameterisations of pion form factors including inelastic effects

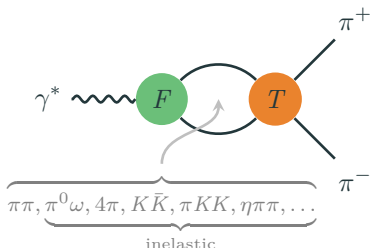
9th Plenary Workshop of the Muon  $g - 2$  Theory Initiative

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**Leon A. Heuser**, G. Chanturia, C. Hanhart, M. Hoferichter, B. Kubis

# Motivation: premise

- Tension in the  $2\pi$  data  
 $\Rightarrow$  no prescription for data-driven  $a_\mu^{\text{HVP, LO}}$  [Aliberti et al., 2025](#)
- Differences in datasets are almost exclusively contained in inelastic contributions [Stoffer et al., 2023](#)
- Inelastic effects are combined as a multiplicative factor in the current dispersive treatment [Leplumey et al., 2025](#)
- No unified fits available for multi-channel treatments  
[N. N. Achasov et al., 2013](#)



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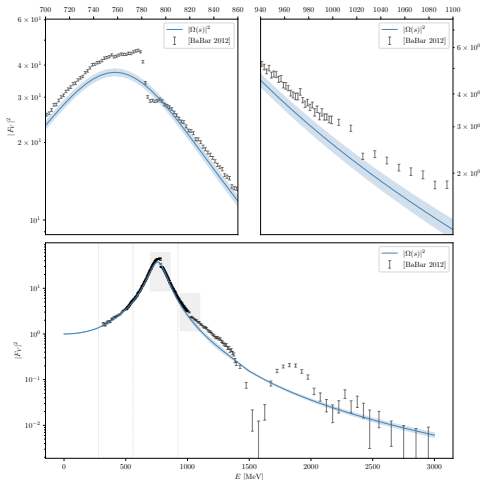
The goal:

- Model 2-pion production + inelastic contributions
- Fit to exclusive data from multiple channels simultaneously

# Framework

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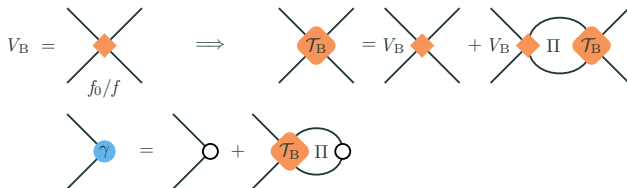
# Two-potential model: why only resonances are not enough



- Pole with channel enough for main features
- Precision not high enough
- Non-resonant contributions improve description
- Unitarity and analyticity necessary

Chanturia, 2022

# Two-potential model: background interactions



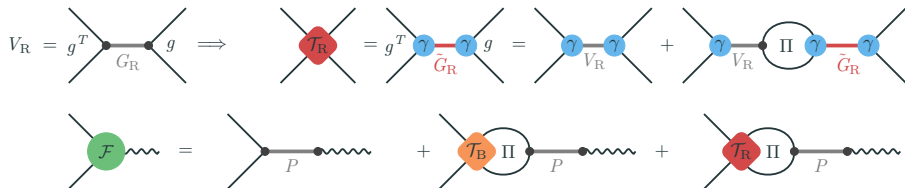
$$V_B(s)_{ab} = \left( \frac{f_0}{f(s)} \right)_{ab} \Rightarrow \mathcal{T}_B = (1 - V_B \Pi)^{-1} V_B$$

Background  $T$ -matrix

$$\text{with } \Pi(s) = \frac{s}{\pi} \int_{s_{\text{thr}}}^{\infty} \frac{ds'}{s'} \frac{\rho(s') \xi(s')^2}{s' - s - i\epsilon} + \text{const.}$$

$$\gamma(s) = 1 + \mathcal{T}_B(s) \Pi(s) \hat{=} p(s) \exp \left( \frac{s}{\pi} \int_{s_{\text{thr}}}^{\infty} \frac{ds'}{s'} \frac{\delta_B(s')}{s' - s - i\epsilon} \right)$$

# Two-potential model: resonance potential



$$V_R(s)_{ab} = \sum_{k=1}^{n_R} \frac{g_a^k g_b^k}{m_k^2 - s} \quad \Rightarrow \quad \boxed{\text{Resonance } T\text{-matrix}} \quad T_R = \gamma^\dagger (1 - V_R \Sigma)^{-1} V_R \gamma$$

with  $\Sigma(s) = \frac{s}{\pi} \int_{s_{\text{thr}}}^{\infty} \frac{ds'}{s'} \frac{\rho(s') \xi(s')^2 |\gamma(s')|^2}{s' - s - i\epsilon} + \text{const.}$

**Form factor**

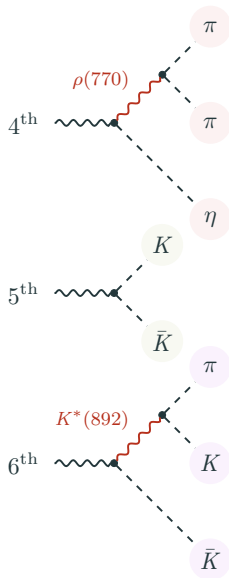
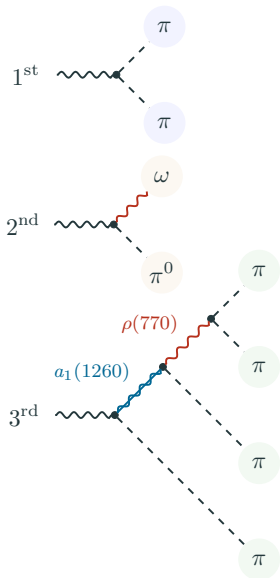
$$\boxed{\mathcal{F} = \gamma (1 - V_R \Sigma)^{-1} P}$$

with  $P(s)_b = \sum_{k=1}^{n_R} \frac{g_b^k a_k}{m_k^2 - s}$

## Application and results

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# Application: open channels ( $J = 1, I = 1$ )



# Application: data

$\pi^+\pi^-$  production: BaBar [Lees et al., 2012b],  
KLOE [Anastasi et al., 2018], CMD-3 [Ignatov et al., 2024]

$\pi\pi$  p-wave phase shift: Bern [Colangelo, 2004, Stoffer et al., 2023, ...]

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- Pion VFF sensitive to phase shifts,  $\omega$  parameters
  - ⇒ individual input for different datasets

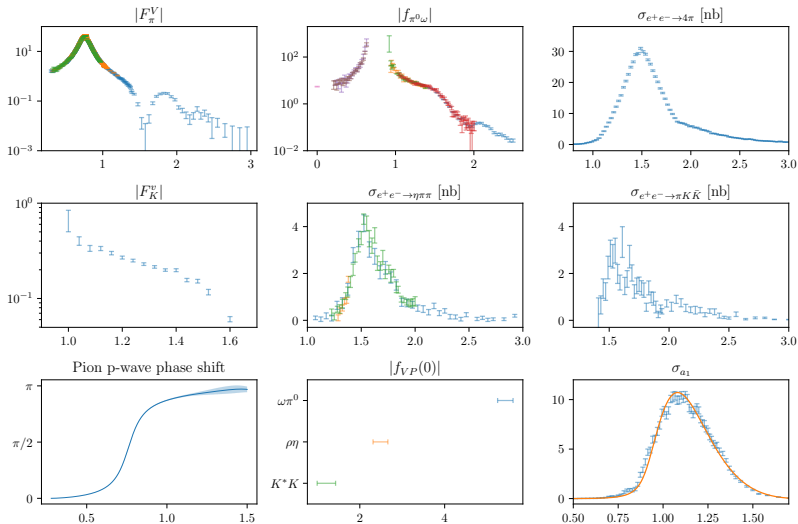
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- Large number of parameters
  - ⇒ must be careful with local minima
  - ⇒ both gradient-based and global minimization routines
- Risk of overfitting
  - ⇒ lowest number of resonances ( $\rho$ ,  $\rho'$ ,  $\rho''$ )
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- Narrow (unphysical) resonances occur as fit artefacts
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- Data quality not uniform
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- Pion VFF sensitive to phase shifts,  $\omega$  parameters
  - ⇒ individual input for different datasets
- Need to correct for bias
  - ⇒ d'Agostini procedure (see e.g. [Ball et al., 2010](#))

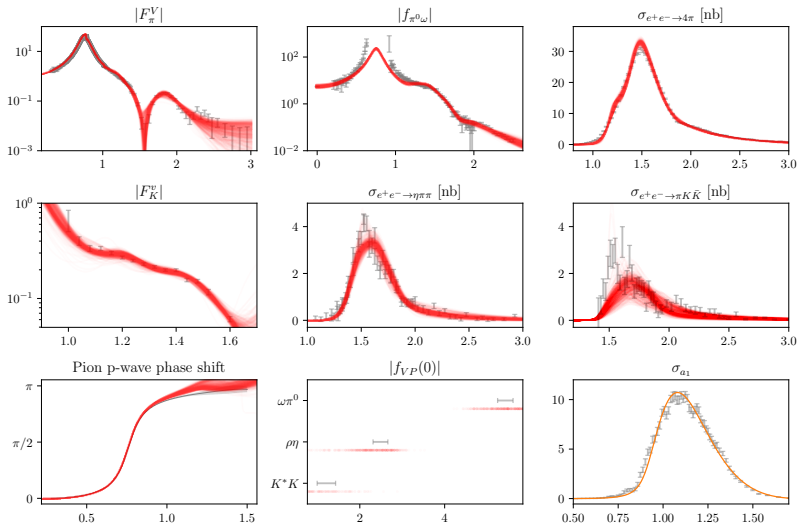
## Application: different fit configurations

| Fit    | Background      | Weights        | Cov. matrices     |             |
|--------|-----------------|----------------|-------------------|-------------|
| Fit #1 | full            | equal          | decorrelated      | next slides |
| Fit #2 | <u>diagonal</u> | equal          | decorrelated      | backup      |
| Fit #3 | full            | <u>unequal</u> | decorrelated      | backup      |
| Fit #4 | full            | equal          | <u>correlated</u> | backup      |

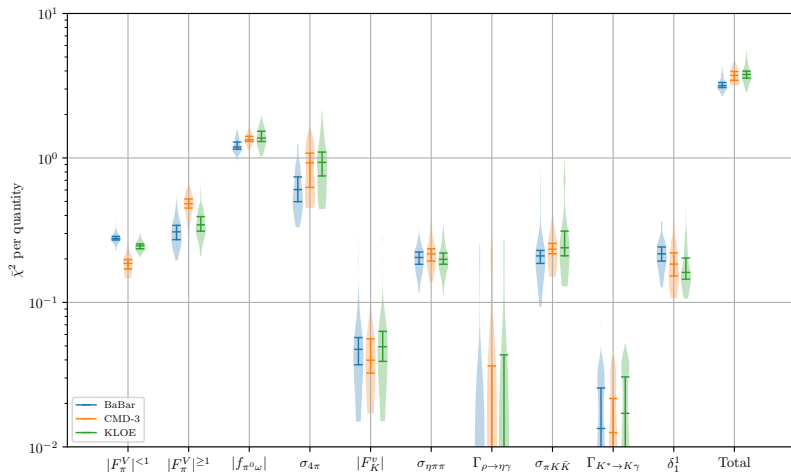
# Application: data



# Results: model performance



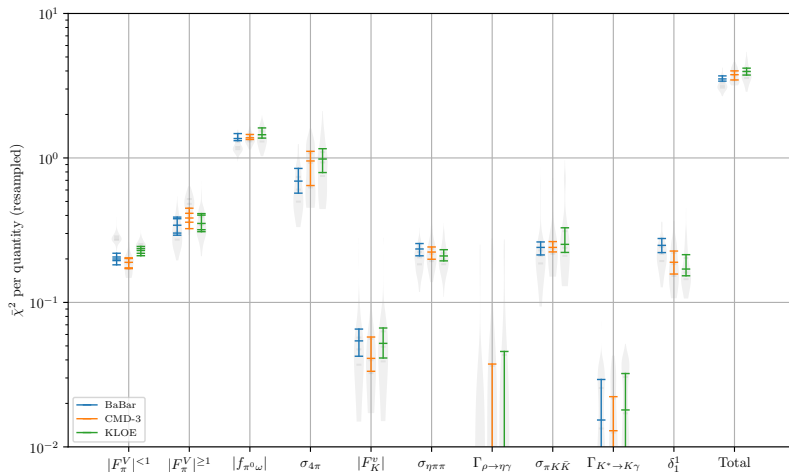
# Results: contributions to the global $\bar{\chi}^2$ (fit #1)



Caution

The  $\pi\pi$  data sets have different number of points

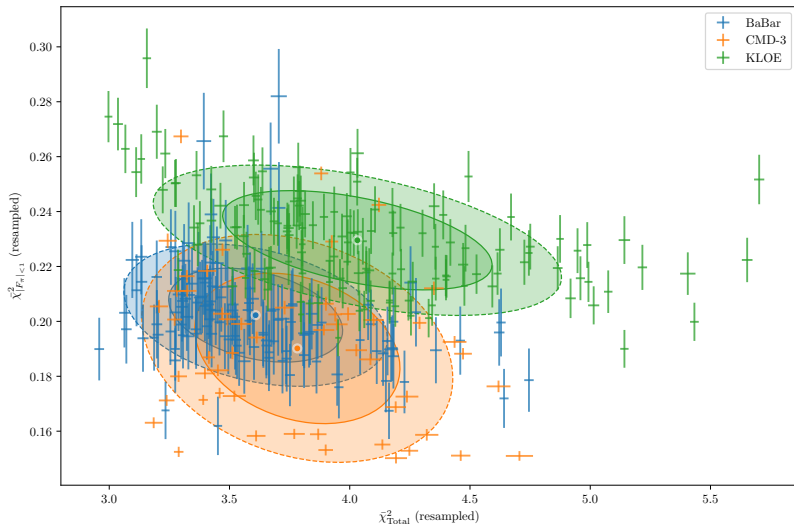
# Results: resampled contributions to $\bar{\chi}^2$ (fit #1)



## Solution

Sample same amount of points for each  $\pi\pi$  data and calculate  $\chi^2$

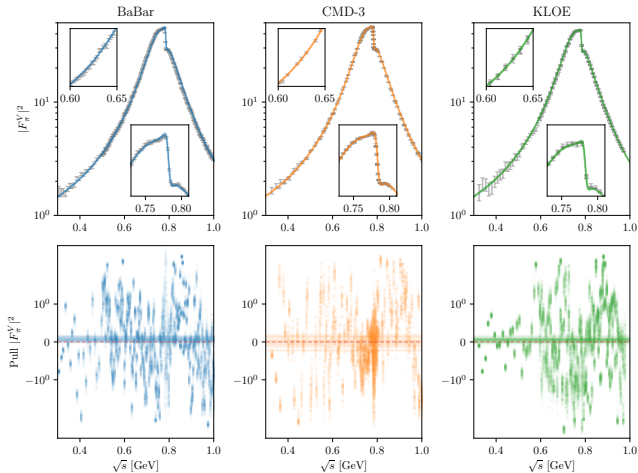
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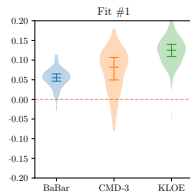
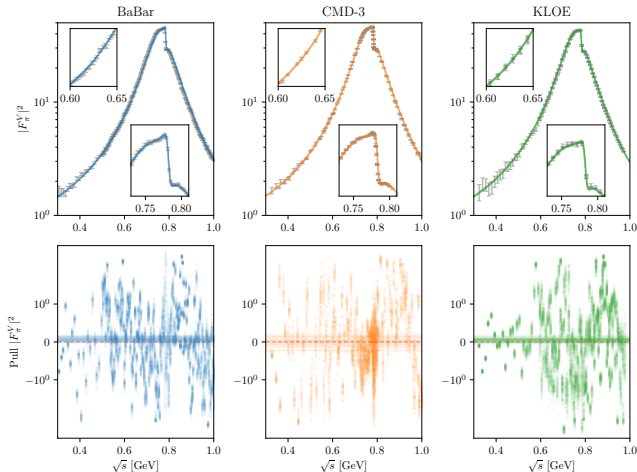
Not enough for proper quantification; actual d.o.f. unknown

# Results: $\pi\pi$ pulls below 1 GeV (fit #1)



$$\text{Pull}(F_i) = (F_{\text{fit},i} - F_{\text{data},i}) / \Delta F_i$$

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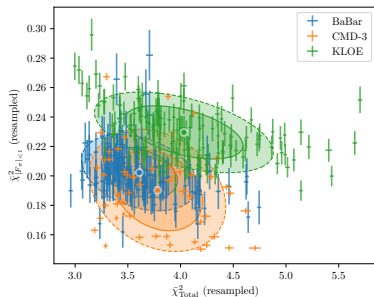
$$\frac{1}{N} \sum_{i=1}^N \text{Pull}(F_i)$$

$$\text{Pull}(F_i) = (F_{\text{fit},i} - F_{\text{data},i}) / \Delta F_i$$

# Results: quantification of differences

- Need no. of effective degrees of freedom of each  $|F_\pi^V| < 1$ 
  - $\implies$  use bootstrapping procedure
  - $\implies$  generate pseudo data around fit and refit
- Peak of  $\chi^2$  distribution is at  $N_{\text{eff}}$
- Use fit from center of ellipses

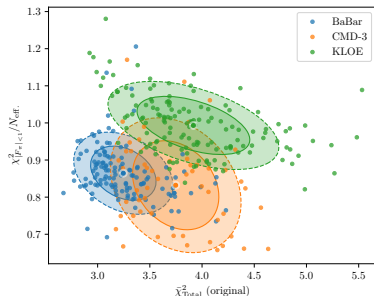
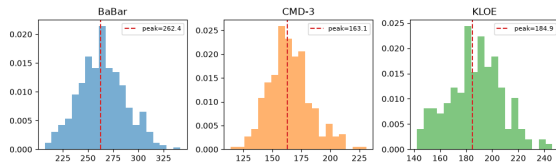
$$N_{\text{BaBar}} = 270 \quad N_{\text{CMD-3}} = 172 \quad N_{\text{KLOE}} = 193$$



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$\implies$   $p$ -value/significance calculation to be done

# Conclusions

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# Summary

- Global fits performed to data from multiple channels
  - Dispersive constraints imposed through  $\pi\pi$ -specific phases
  - Two-potential formalism ensures unitarity and analyticity
  - To be submitted to arXiv soon!
- 

## Qualitative:

- Differences between BaBar and CMD-3 small, KLOE noticeably worse
  - All fits have positive pulls
- 

## Quantitative:

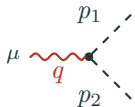
- Proper significance calculation to be done

Thank you for your  
attention!

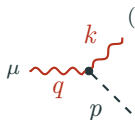
## Backup slides

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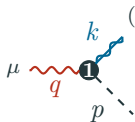
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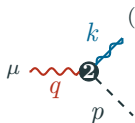
$$\sim \xi_{VPP}^{\mu} = (p_1 - p_2)^{\mu} \quad (\text{e.g. } \gamma^* \rightarrow \pi\pi, K\bar{K})$$



$$\sim \xi_{VVP}^{(\lambda)\mu} = \epsilon^{\mu\nu\alpha\beta} n_{\nu}^{(\lambda)} p_{\alpha} q_{\beta} \quad (\text{e.g. } \gamma^* \rightarrow \omega\pi^0, \rho\eta, K^*K)$$



$$\sim \xi_{VAP,1}^{(\lambda)\mu} = (q \cdot k) n^{(\lambda)\mu} - (n^{(\lambda)} \cdot q) k^{\mu} \quad (\text{e.g. } \gamma^* \rightarrow a_1\pi)$$



$$\sim \xi_{VAP,2}^{(\lambda)\mu} = q^2 k^2 n^{(\lambda)\mu} + (q \cdot k)(n^{(\lambda)} \cdot q) q^{\mu} - k^2 (n^{(\lambda)} \cdot q) q^{\mu}$$

# Application: semistable channels

Spectral function for an unstable state  $A$  via dressed propagator  
→ fix parameters  $g$ ,  $M_A$  like in [Heuser et al., 2024](#) via pole position

$$\sigma_A(q^2) = -\frac{1}{\pi} \text{Im} (G_A(q^2)) , \quad \text{where} \quad G_A(q^2) = \frac{1}{q^2 - M_A^2 + g^2 \Pi_i(q^2)}$$

---

Imaginary part of the self-energy  $\Pi_{AB}$  with  $A$  being a (broad) resonance

$$\text{Im}\Pi_{AB} = \int_{s_{\text{thr}, \text{decay}}}^{(\sqrt{s} - M_B)^2} \text{d}q_A^2 \rho(s, q_A^2, M_B^2) \xi^2(s, q_A^2, M_B^2) \sigma_A(q_A^2)$$

---

$$\text{Im}\Pi_{\pi\pi\eta}(s) = \int_{(2M_\pi)^2}^{(\sqrt{s} - M_\eta)^2} \text{d}q_\rho^2 \sigma_\rho(q_\rho^2) \rho(s, q_\rho^2, M_\eta^2) \xi_{VVP}^2(s, q_\rho^2, M_\eta^2)$$

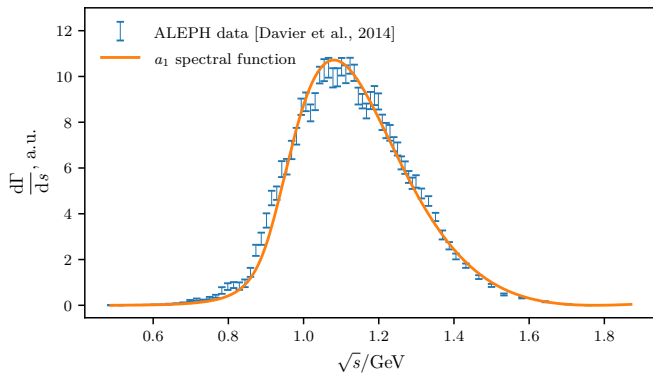
$$\text{Im}\Pi_{\pi KK}(s) = \int_{(M_\pi + M_K)^2}^{(\sqrt{s} - M_K)^2} \text{d}q_{K^*}^2 \sigma_{K^*}(q_{K^*}^2) \rho(s, q_{K^*}^2, M_K^2) \xi_{VVP}^2(s, q_{K^*}^2, M_K^2)$$

# Application: $4\pi$ and the $a_1$ spectral function

$a_1\pi$  problematic:

$a_1$  pole might be model dependent + three-body effects

→ use  $\tau$ -decay data for  $a_1$  and fit  $g_{a_1\rho\pi}$ ,  $M_{a_1}$



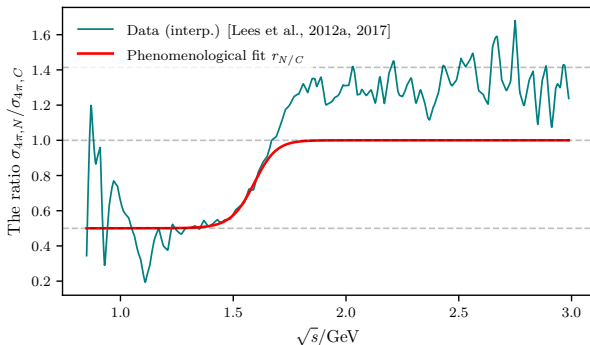
# Application: $4\pi$ and the $a_1$ spectral function

$a_1\pi$  problematic:

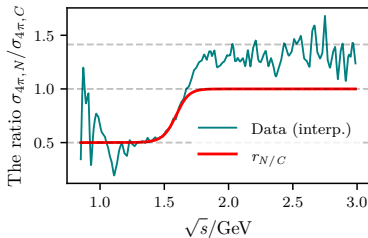
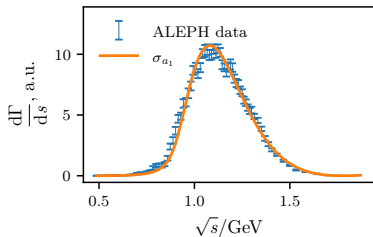
difference between interference effects in  $2\pi^0\pi^+\pi^-$  and  $2(\pi^+\pi^-)$  not accounted for within the model

→ fit multiplicative factor to ratio of corresponding cross sections

$$r_{N/C}(s) = \frac{1}{2} \left( 1 + \exp^{-1} \left( \frac{a^2 - s}{b^2} \right) \right)$$



# Application: $4\pi$ and the $a_1$ spectral function



$$\text{Im}\Pi_{\rho\pi}(q_{a_1}^2) = \int_{(2M_\pi)^2}^{(\sqrt{q_{a_1}^2} - M_\pi)^2} dq_\rho^2 \sigma_\rho(q_\rho^2) \rho(q_{a_1}^2, q_\rho^2, M_\pi^2) \xi_{VVP}^2(q_{a_1}^2, q_\rho^2, M_\pi^2)$$

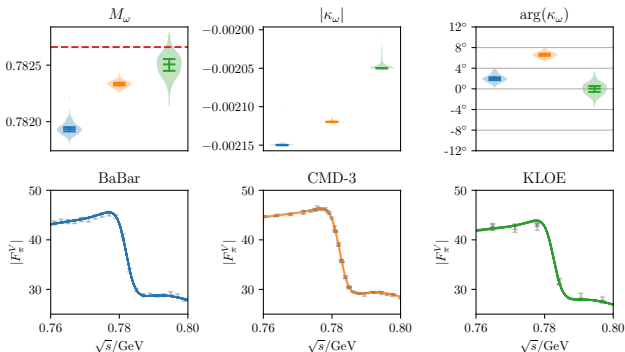
$$\sigma_{a_1}(q_{a_1}^2) = -\frac{1}{\pi} \text{Im} \left( \frac{1}{q_{a_1}^2 - M_{a_1}^2 + g_{a_1\rho\pi}^2 \Pi_{\rho\pi}(q_{a_1}^2)} \right)$$

$$\text{Im}\Pi_{i,4\pi}(s) = \int_{(3M_\pi)^2}^{(\sqrt{s} - M_\pi)^2} dq_{a_1}^2 \sigma_{a_1}(q_{a_1}^2) \rho(s, q_{a_1}^2, M_\pi^2) \xi_{i,VAP}^2(s, q_{a_1}^2, M_\pi^2) r_{N/C}(s)$$

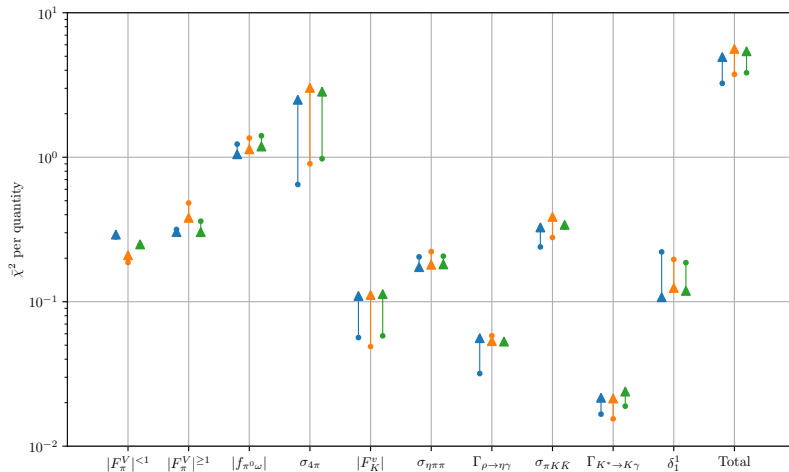
# Results: $\omega$ parameters from 2-pion data

$$F_{\pi}^V(s) = \left(1 + \kappa_{\omega} \frac{s}{s - M_{\omega}^2 + iM_{\omega}\Gamma_{\omega}} + \dots\right) F_{\pi}^{V,I=1}(s)$$

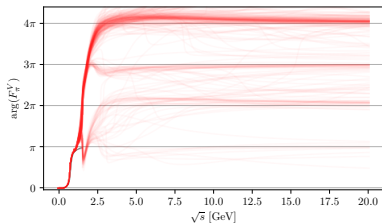
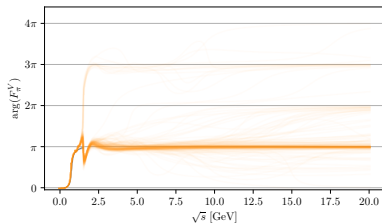
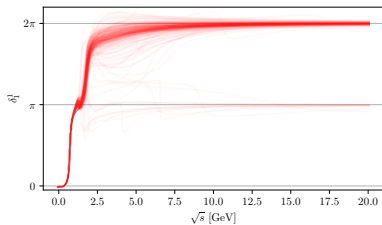
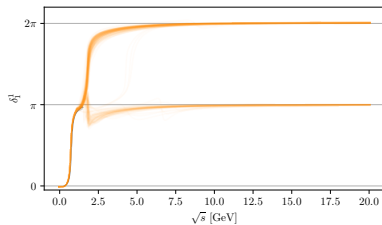
- $\Gamma_{\omega}$  fixed from  $3\pi$  [Hoferichter et al., 2023](#)
- Bounds for  $M_{\omega}$  and  $\kappa_{\omega}$  were selected according to the  $\pi\pi$  dataset [\[Stoffer et al., 2023\]](#)



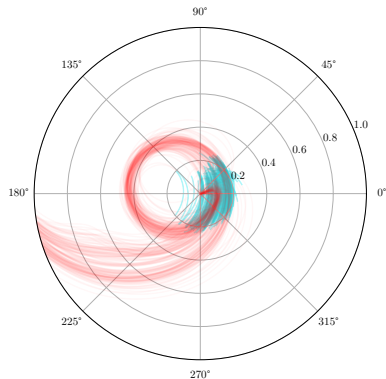
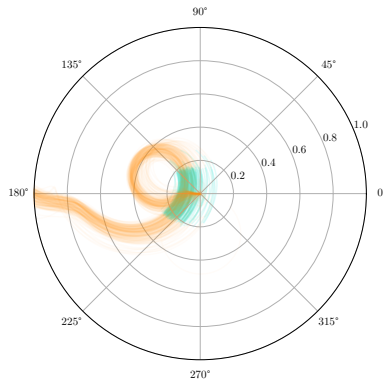
# Results: fit #2 (diagonal background)



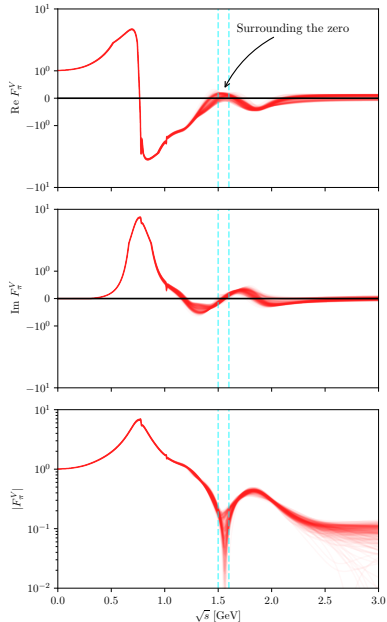
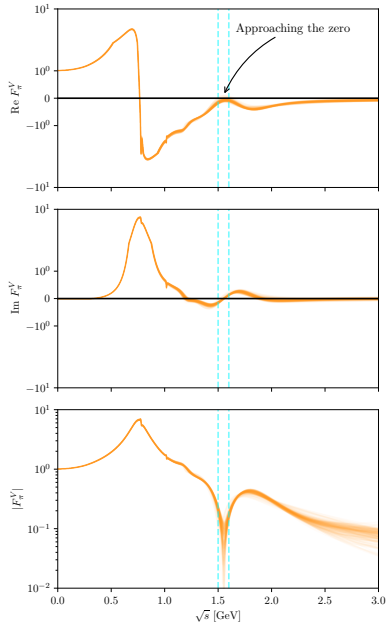
# Results: fit #2 (diagonal background)



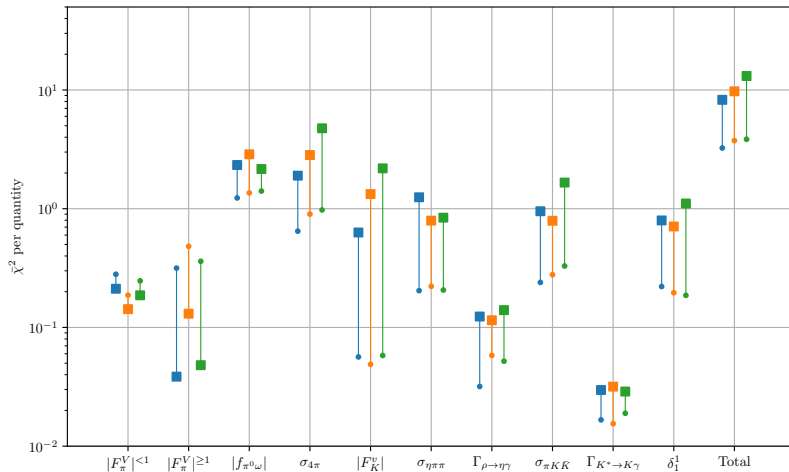
## Results: fit #2 (diagonal background)



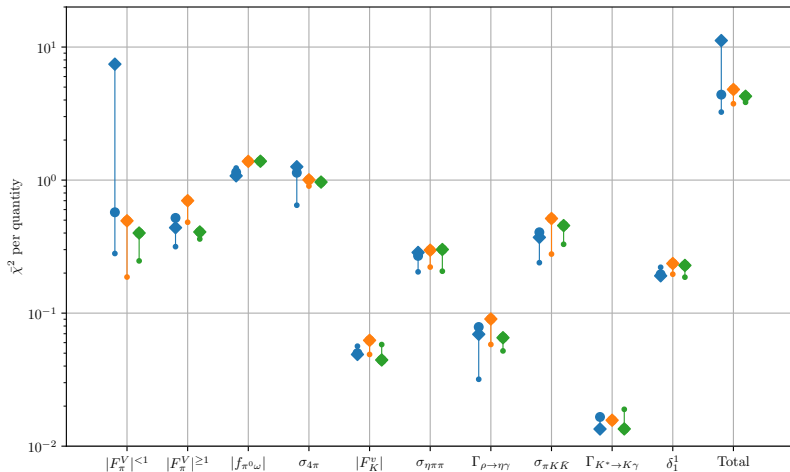
# Results: fit #2 (diagonal background)



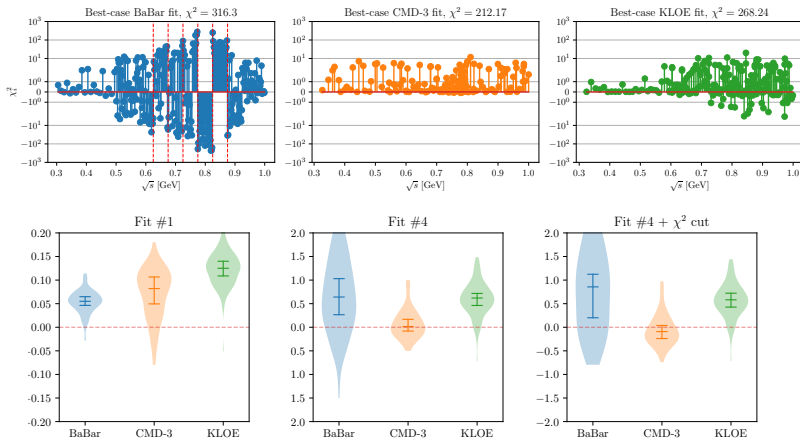
# Results: fit #3 (exaggerated $\pi\pi$ )



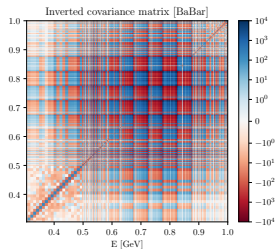
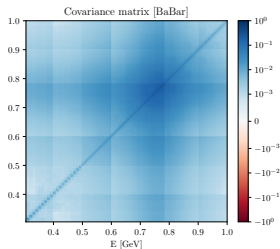
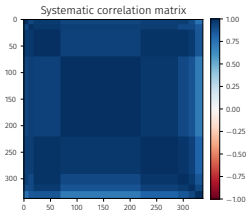
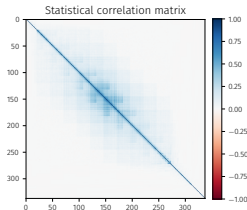
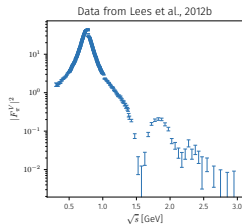
# Results: fit #4 (correlation included)



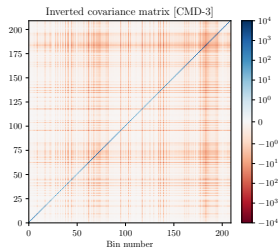
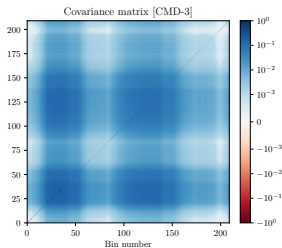
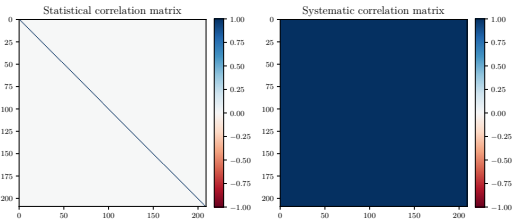
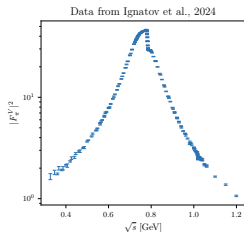
# Results: fit #4 (correlation included)



# Covariance matrices: BaBar

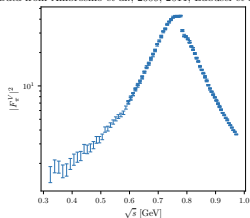


# Covariance matrices: CMD-3

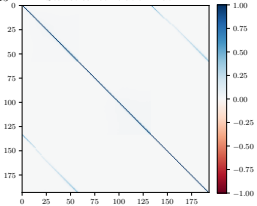


# Covariance matrices: KLOE

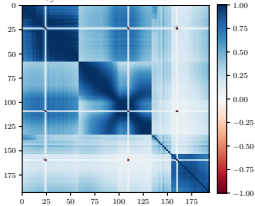
Data from Ambrosino et al., 2009, 2011; Babusci et al., 2013



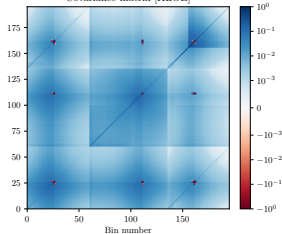
Statistical correlation matrix



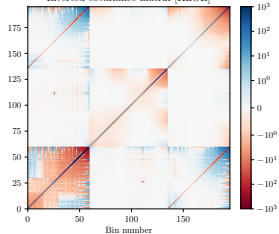
Systematic correlation matrix



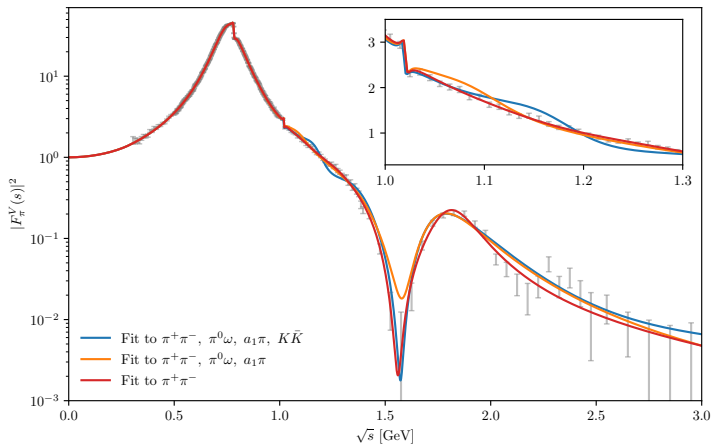
Covariance matrix [KLOE]



Inverted covariance matrix [KLOE]



# Emergent structure after the $\phi$



# Literature

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